

A convergence result for spatially inhomogeneous Preisach operators

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Abstract. Based on the properties of the strongly nonlinear Preisach operator, a convergence result is proved under reasonable assumptions.

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Keywords.

1. Introduction

It is natural to consider materials whose properties are spatially inhomogeneous, i.e. differ at different points.

In this paper we will consider spatially inhomogeneous Preisach hysteresis operators and give a result on continuity of Preisach operators with respect to convergence of density functions.

Hysteresis is a nonlinear phenomena. The basic feature of hysteresis behaviour is a memory effect and irreversibility of the process. A systematic mathematical investigation of hysteresis operators started relatively recently, see e.g. [1], [2], [7], [8], although e.g. the Preisach model itself was introduced much earlier [9].

Our result will be critical for a homogenization argument for a parabolic equation with Preisach hysteresis, analogous to Proposition 2.12 in [6]. This will be published elsewhere.

In [3] a corresponding convergence result is established for the Prandtl-Ishlinskii hysteresis operator, whose nonlinearity is simpler than for the Preisach operator considered here.

The paper is organized as follows. Section 1 is devoted to a description of a simple hysteresis operator, the play, and its basic properties. The Preisach operator is defined in Section 2 and we list also there its important properties. The main result is the subject of Section 3.

2. The Play operator

Let $r > 0$ be a parameter, $u(t)$ a continuous input function on the time interval $I = [0, T]$ and $w_r^0 \in [-r, r]$ an initial state. We consider a variational inequality

$$\mathcal{G}(t) \in [-r, r], t \in I \quad (1)$$

$$\left(\dot{\mathcal{G}}(t) - \dot{u}(t)\right) (\phi - \mathcal{G}(t)) \geq 0 \text{ for a.e. } t \in I, \text{ for all } \phi \in [-r, r] \quad (2)$$

$$\mathcal{G}(0) = w_r^0 \quad (3)$$

for the unknown $\mathcal{G}(t)$. For an input $u \in W^{1,1}(I)$ this problem admits a unique solution $\mathcal{G}_r[u, w_r^0] \in W^{1,1}(I)$. The play operator \mathcal{E}_r with threshold r is defined by the relation

$$\mathcal{E}_r[u, w_r^0] = u - \mathcal{G}_r[u, w_r^0], \quad (4)$$

see Figure 1.

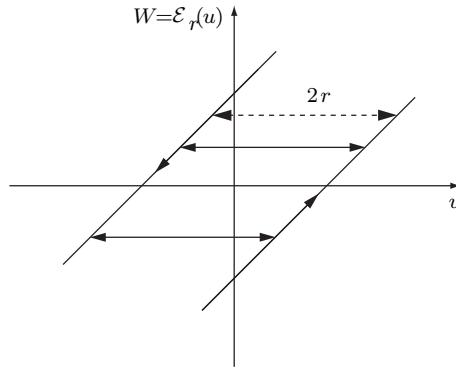


Figure 1

For piecewise monotone inputs, in each interval of monotonicity $[t_0, t_1]$ of the input function $u(t)$, the relation

$$\mathcal{E}_r[u, w_r^0](t) = \max \{u(t) - r, \min\{u(t) + r, \mathcal{E}_r[u, w_r^0](t_0)\}\} \quad (5)$$

follows from (1) and is often used for an alternative definition of the play operator, see e.g. [8]. In the following we use the “unperturbed” or “virgin” initial state defined by

$$w_r^0 = \min \{r, \max\{-r, u(0)\}\}. \quad (6)$$

and write only $\mathcal{E}_r[u]$.

The play operator has also a geometric “piston-in-cylinder” interpretation. Consider a cylinder of length $2r$ with a piston moving inside. If the input $u(t)$

denotes the position of the piston moving in the cylinder, then the position of the cylinder yields the play operator.

We survey important properties of the play operator.

Theorem 1. (i) (see e.g. Proposition II.1.1 in [7]) The play operator $\mathcal{E}_r[u]$ is Lipschitz continuous in $W^{1,1}(0, T)$.

(ii) (see e.g. Proposition II.4.5 in [7]) The play operator admits a Lipschitz continuous extension to $C([0, T])$ and

$$|\mathcal{E}_r[u_1](t) - \mathcal{E}_r[u_2](t)| \leq \max_{0 \leq s \leq t} |u_1(s) - u_2(s)|, \forall t \in I. \tag{7}$$

(iii) (see e.g. Lemma 2.3.8 in [1]) The play operator is piecewise monotone in the following sense:

$$\frac{d\mathcal{E}_r[u]}{dt} \cdot \frac{du}{dt} \geq 0 \text{ a.e. in } (0, T). \tag{8}$$

3. The Preisach operator

The simplest example of a hysteresis nonlinearity is given by a switch or relay with hysteresis, $h_{v,r} : C([0, T]) \times \{-1, 1\} \rightarrow BV(0, T)$ with input u (magnetic field) and output $h_{v,r}$ (magnetization), see Figure 2. The relay is characterized by two parameters $v \in R^1$ (interaction field) and $r > 0$ (critical field of coercivity) and is defined formally as follows: Let R_+^2 denote the set $\{(v, r) \in R^2; r > 0\}$. For given parameters $(v, r) \in R_+^2$, input $u \in C([0, T])$, initial magnetization $\xi \in \{-1, 1\}$ and any time $t \in [0, T]$, put

$$X_t := \{\tau \in (0, T]; |u(\tau) - v| = r\}. \tag{9}$$

We then define

$$h_{v,r}[u, \xi](0) = \begin{cases} -1 & \text{if } u(0) \leq v - r \\ \xi & \text{if } v - r < u(0) < v + r \\ 1 & \text{if } u(0) \geq v + r \end{cases} \tag{10}$$

and

$$h_{v,r}[u, \xi](t) = \begin{cases} h_{v,r}[u, \xi](0) & \text{if } X_t = \emptyset \\ -1 & \text{if } X_t \neq \emptyset \text{ and } u(\max X_t) = v - r \\ 1 & \text{if } X_t \neq \emptyset \text{ and } u(\max X_t) = v + r, \end{cases} \tag{11}$$

see Figure 2.

It is often convenient to use the following representation of the relay by means of the system $\mathcal{E}_r, r > 0$ of play operators.

Lemma 2. (See e.g. Lemma II.3.6 in Krejčí for more general initial states). Let $u \in C([0, T])$ be given. For $(v, r) \in R_+^2$, put $\xi := -1$ if $v \geq 0$, $\xi = 1$ if $v < 0$.

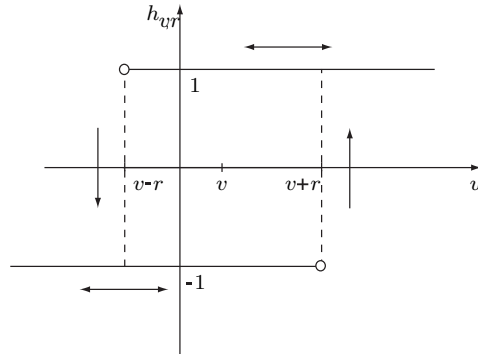


Figure 2

Then for every $t \in [0, T]$ and $(v, r) \in \mathbb{R}_+^2$, $v \neq \mathcal{E}_r[u](t)$ we have

$$h_{v,r}[u, \xi](t) = \begin{cases} -1 & \text{if } v \geq \mathcal{E}_r[u](t) \\ 1 & \text{if } v \leq \mathcal{E}_r[u](t). \end{cases} \quad (12)$$

The output of the Preisach model is formally defined as an average over all elementary switches with a given density function $\psi \in L_{loc}^1(\mathbb{R}_+^2)$ by the formula (see e.g. Krasnosel'skii and Pokrovskii [2])

$$\mathcal{P}[u](t) := \int_0^\infty \int_{-\infty}^\infty \psi(v, r) h_{v,r}[u, \xi](t) dv dr, \quad (13)$$

where the initial values of the relays are taken as -1 if $v > 0$ and $+1$ otherwise. Using Lemma 2 on the representation of the relay by a system of plays the output of the Preisach operator can be expressed as

$$\mathcal{P}[u](t) = C + \int_0^\infty g(\mathcal{E}_r[u](t), r) dr, \quad (14)$$

where

$$g(v, r) = \int_0^v \psi(z, r) dz, \quad (15)$$

C is a constant and $\mathcal{E}_r[u](t)$ denotes the play operator.

Remark 1. Notice that the integral in (14) is meaningful if $u \in C([0, T])$ since $\mathcal{E}_r[u](t) = 0$ for r sufficiently large and $g(0, r) = 0$ for all $r > 0$.

In the sequel we will use the following assumptions :

(P1) There exists $\beta \in L_{loc}^1(0, \infty)$, $\beta \geq 0$ a.e. such that

$$0 \leq \psi(z, r) \leq \beta(r) \text{ for a.e. } (z, r) \in \mathbb{R}_+^2. \quad (16)$$

For $R > 0$ put $b(R) := \int_0^R \beta(r) dr$.

(P2)

$$\frac{\partial \psi}{\partial z} \in L_{loc}^{\infty}(R_+^2). \quad (17)$$

The next theorem shows conditions under which the Preisach operator is Lipschitz continuous on $C^0([0, T])$. Proof of the extended version of this theorem (for more general initial inputs) can be found in [7], Proposition II.3.11.

Theorem 3. *Let the Assumption (P1) be satisfied and let $R > 0$ be given. Then for every $u, v \in C([0, T])$ such that $\|u\|_{C^0([0, T])}, \|v\|_{C^0([0, T])} \leq R$ the Preisach operator (14) satisfies*

$$\|\mathcal{P}[u] - \mathcal{P}[v]\|_{C^0([0, T])} \leq b(R)\|u - v\|_{C^0([0, T])}. \quad (18)$$

Lemma 4. *(Lemma II.4.1 in [7]) Let the Assumptions (P1) and (P2) be satisfied. Then for $u \in W^{1,1}(0, T)$, $r > 0$ and $t \in [0, T]$, we have $\mathcal{P} \in W^{1,1}(0, T)$ and for a.e. $t \in [0, T]$*

$$\dot{\mathcal{P}}[u](t) = \int_0^{\infty} \dot{\mathcal{E}}_r[u](t) \psi(\mathcal{E}_r[u](t), r) dr. \quad (19)$$

It follows from the previous lemma and from the definition of the play operator that the Preisach operator is piecewise monotone. We have

Theorem 5. *The Preisach operator is under the Assumptions (P1) and (P2) piecewise monotone, this means that for $u \in W^{1,1}(0, T)$*

$$\dot{\mathcal{P}}[u](t) \dot{u}(t) \geq 0 \text{ a.e.} \quad (20)$$

We listed only some properties of the hysteresis operators which will be useful for our purposes. For a full exposition of hysteresis operators and their properties we recommend [1], [2], [7], [8]. In each of these an interested reader can find a different approach.

4. Spatially dependent Preisach operators

We will consider the spatially dependent constitutive relation described by the Preisach operator with a spatially dependent density function $\psi(x, z, r) \in L_{loc}^1(\Omega \times R_+^2)$, $\Omega \subset R^n$ is considered to be a bounded domain with Lipschitz boundary.

For a given input $u(x, t) : \Omega \times I \rightarrow R$ we define the output of the spatially dependent Preisach operator as follows:

$$\mathcal{P}[u](x, t) = C + \int_0^{\infty} g(\mathcal{E}_r[u(x, \cdot)](t), r) dr, \quad (21)$$

where

$$g(x, v, r) = \int_0^v \psi(x, z, r) dz. \quad (22)$$

Theorem 6. Let $\psi_n \in L^p(\Omega, L^1_{loc}(R_+^2))$, $p \geq 1$, be a sequence of space dependent density functions satisfying the Assumption (P1) for a.e. $x \in \Omega$. Assume that ψ_n converge to ψ in $L^p(\Omega, L^1_{loc}(R_+^2))$. Let $\mathcal{P}_n, \mathcal{P}$ be the Preisach operators corresponding to ψ_n, ψ respectively. Let u_n be a sequence in $L^p(\Omega, C(I))$ such that $u_n(x, \cdot) \in C(I)$ for a.e. $x \in \Omega$ and $\|u_n - u\|_{C(I)} \rightarrow 0$ as $n \rightarrow \infty$.

Then $\mathcal{P}_n[u_n](\cdot, t)$ converge to $\mathcal{P}[u](\cdot, t)$ for every $t \in I$ in $L^p(\Omega)$.

Proof. We have for a.e. $x \in \Omega$ and every $t \in I$

$$\int_0^\infty \int_0^{\mathcal{E}_r[u_n](t)} \psi_n(x, z, r) dz dr - \int_0^\infty \int_0^{\mathcal{E}_r[u](t)} \psi(x, z, r) dz dr = \quad (23)$$

$$= \int_0^\infty \int_0^{\mathcal{E}_r[u_n](t)} \psi_n(x, z, r) dz dr - \int_0^\infty \int_0^{\mathcal{E}_r[u](t)} \psi_n(x, z, r) dz dr + \quad (24)$$

$$+ \int_0^\infty \int_0^{\mathcal{E}_r[u](t)} [\psi_n(x, z, r) - \psi(x, z, r)] dz dr = \quad (25)$$

$$\int_0^\infty \int_{\mathcal{E}_r[u](t)}^{\mathcal{E}_r[u_n](t)} \psi_n(x, z, r) dz dr + \int_0^\infty \int_0^{\mathcal{E}_r[u](t)} [\psi_n(x, z, r) - \psi(x, z, r)] dz dr. \quad (26)$$

The first integral on the right hand side of the last expression can be estimated using the Assumption (P1) as follows

$$\int_0^\infty \int_{\mathcal{E}_r[u](t)}^{\mathcal{E}_r[u_n](t)} \psi_n(x, z, r) dz dr \leq \int_0^\infty \int_{\mathcal{E}_r[u](t)}^{\mathcal{E}_r[u_n](t)} \beta(r) dz dr = \quad (27)$$

$$= \int_0^\infty \beta(r) [\mathcal{E}_r[u_n](t) - \mathcal{E}_r[u](t)] dr \leq \int_0^R \beta(r) |\mathcal{E}_r[u_n](t) - \mathcal{E}_r[u](t)| dr, \quad (28)$$

for some $R > 0$, see Remark 1.

The later term can be further estimated using the Lipschitz continuity of the play operator in $C([0, T])$ (Theorem 1 (ii)) as follows:

$$\int_0^R \beta(r) |\mathcal{E}_r[u_n](t) - \mathcal{E}_r[u](t)| dr \leq \|u_n - u\|_{C^0([0, T])} b(R), \quad (29)$$

where $b(R)$ is defined in (P1).

The estimates above imply that for every $t \in I$

$$\|\mathcal{P}_n[u_n](t) - \mathcal{P}[u](t)\|_{L^p(\Omega)} \leq b(R) \|u_n - u\|_{L^p(\Omega, C^0([0, T]))} + \quad (30)$$

$$+ \left\| \int_0^\infty \int_0^{\mathcal{E}_r[u](t)} [\psi_n(x, z, r) - \psi(x, z, r)] dz dr \right\|_{L^p(\Omega)}. \quad (31)$$

The first term on the right hand side of the last inequality converges by assumptions to 0. To estimate the second term, again by Remark 1 $\mathcal{E}_r[u](t) = 0$ for r sufficiently large, and if $u(x, \cdot) \in C([0, T])$ for a.e. $x \in \Omega$ by Theorem 1 (ii)

$\mathcal{E}_r[u](x, \cdot) \in C([0, T])$, so the integral over r is through a finite interval, and converges to zero because of the assumption on the convergence of ψ_n . The statement follows. \square

Remark 2. It follows from the proof of the previous theorem that the convergence ψ_n to ψ in $L^p(\Omega, L^1_{loc}(R^2_+))$ can be easily replaced by the convergence in $L^\infty(\Omega \times R^2_+)$ weakly star, as is typically the case we get in homogenization arguments, getting the weak star convergence of the Preisach operators in $L^\infty(\Omega)$.

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